

# A kinetic neutral atom model for tokamak scrape-off layer turbulence simulations

Christoph Wersal, Paolo Ricci, Federico Halpern, Fabio Riva



CRPP - EPFL

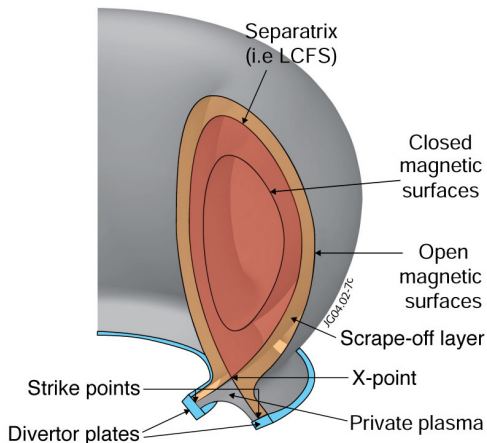
SPS Annual Meeting 2014

02.07.2014

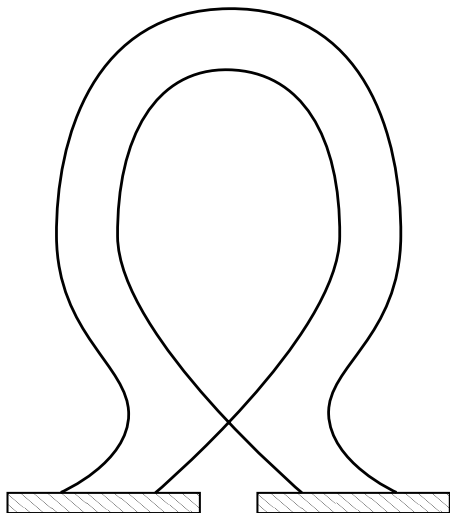


# The tokamak scrape-off layer (SOL)

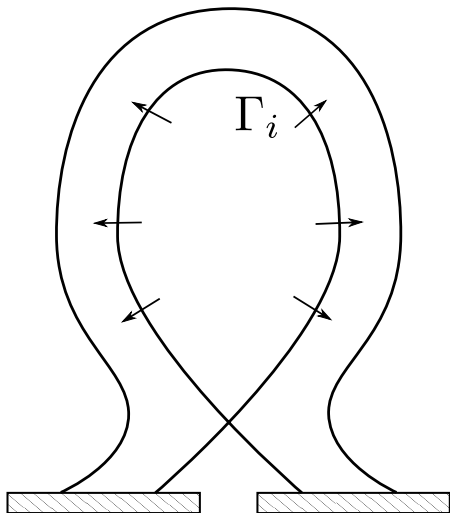
- ▶ Open field lines
- ▶ Heat exhaust
- ▶ Confinement
- ▶ Impurities
- ▶ Fusion ashes removal
- ▶ Fueling the plasma (recycling)
- ▶ Three regimes
  - ▶ Convection limited
  - ▶ Conduction limited
  - ▶ Detached



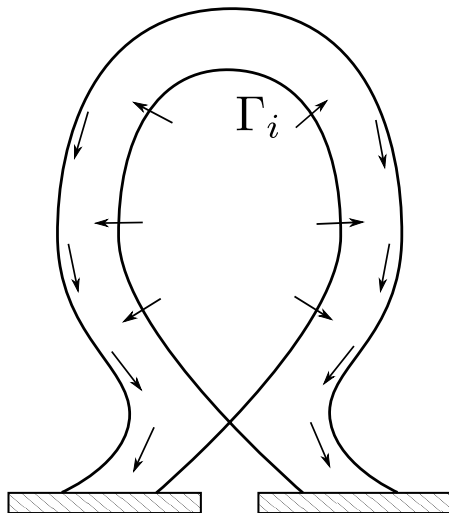
# Convection limited regime



# Convection limited regime

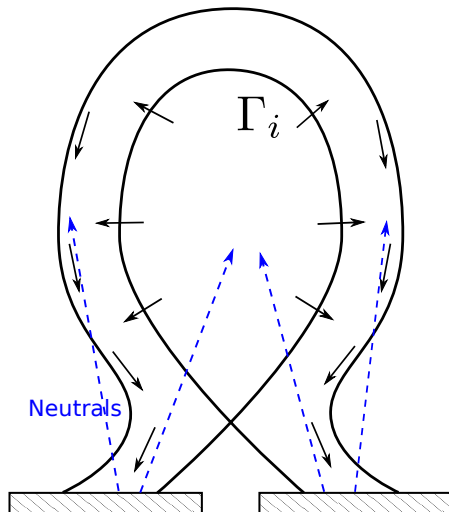


# Convection limited regime



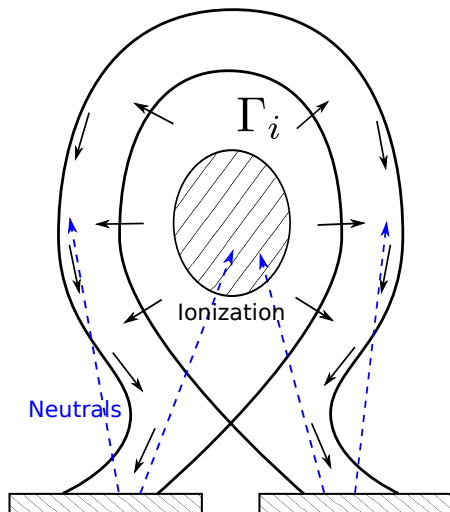
# Convection limited regime

- ▶ Low plasma density
- ▶ Long  $\lambda_{mfp}$  for neutrals



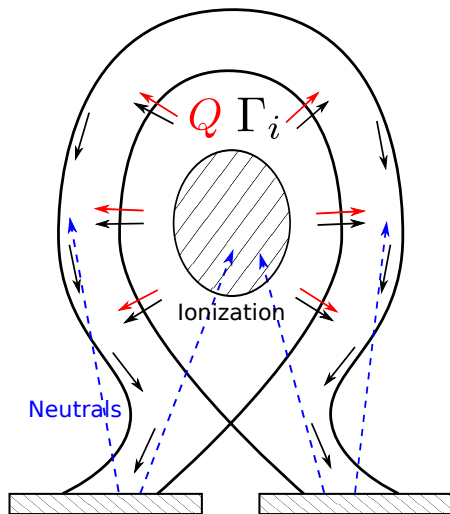
# Convection limited regime

- ▶ Low plasma density
- ▶ Long  $\lambda_{mfp}$  for neutrals
- ▶ Ionization in the core



# Convection limited regime

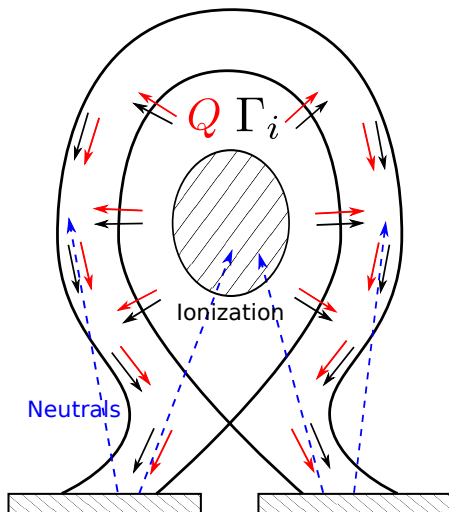
- ▶ Low plasma density
- ▶ Long  $\lambda_{mfp}$  for neutrals
- ▶ Ionization in the core
- ▶ Heat  $\approx$  particle source





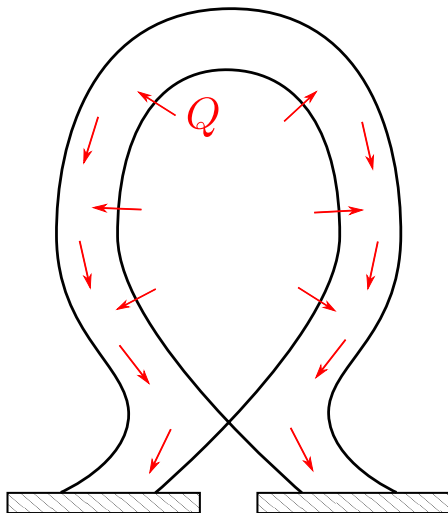
# Convection limited regime

- ▶ Low plasma density
- ▶ Long  $\lambda_{mfp}$  for neutrals
- ▶ Ionization in the core
- ▶ Heat  $\approx$  particle source
- ▶  $Q$  is mainly convective



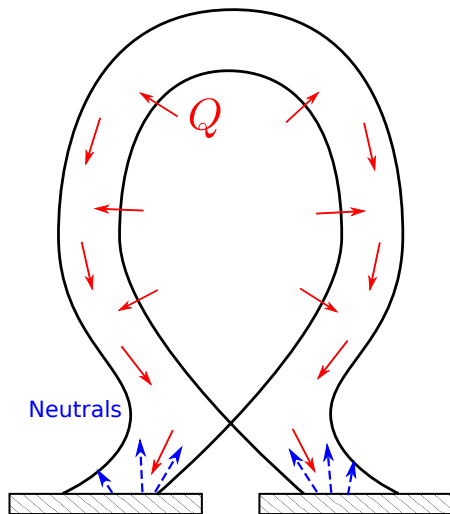
# Conduction limited regime

- ▶ High plasma density



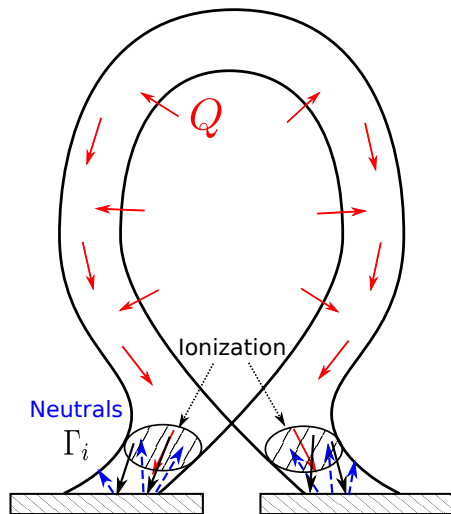
# Conduction limited regime

- ▶ High plasma density
- ▶ Short  $\lambda_{mfp}$



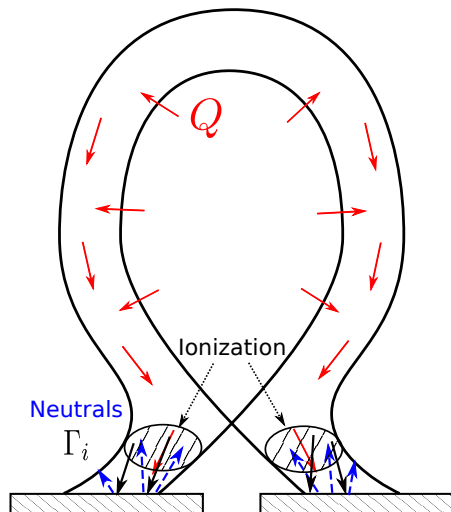
# Conduction limited regime

- ▶ High plasma density
- ▶ Short  $\lambda_{mfp}$
- ▶ Ionization close to targets



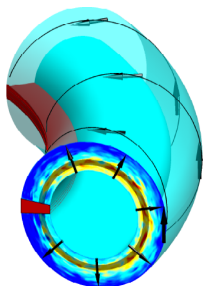
# Conduction limited regime

- ▶ High plasma density
- ▶ Short  $\lambda_{mfp}$
- ▶ Ionization close to targets
- ▶ Temperature gradients form
- ▶  $Q$  is mainly conductive



# The GBS code, a tool to simulate SOL turbulence

- ▶ Drift-reduced Braginskii equations  
 $d/dt \ll \omega_{ci}, k_{\perp}^2 \gg k_{\parallel}^2$
- ▶ Evolves scalar fields in 3D geometry  
 $n, \Omega, v_{\parallel e}, v_{\parallel i}, T_e, T_i$
- ▶ Flux-driven, no separation between equilibrium and fluctuations
- ▶ Power balance between plasma outflow from the core, turbulent transport, and sheath losses [Ricci et al., PPCF, 2012]
- ▶ No divertor geometry
- ▶ No neutral physics



# The model

Kinetic model for the neutral atoms with Krook operators

$$\frac{\partial f_n}{\partial t} + \vec{v} \cdot \frac{\partial f_n}{\partial \vec{x}} = -v_{ion} f_n - v_{cx}(f_n - n_n \Phi_i) + v_{rec} n_i \Phi_i \quad (1)$$

$$v_{ion} = n_e \langle v_e \sigma_{ion}(v_e) \rangle, \quad v_{cx} = n_i \langle v_{rel} \sigma_{cx}(v_{rel}) \rangle$$

$$v_{rec} = n_e \langle v_e \sigma_{rec}(v_e) \rangle, \quad \Phi_i = f_i / n_i$$

Boundary conditions

( $v_{\perp}$  in respect to the surface;  $\theta$  between  $\vec{v}$  and normal vector to the surface)

$$\int d\vec{v} v_{\perp} f_n(\vec{x}_w, \vec{v}) + u_{i\perp} n_i = 0 \quad (2)$$

$$f_n(\vec{x}_w, \vec{v}) \propto \cos(\theta) e^{mv^2/2T_w} \quad \text{for } v_{\perp} > 0$$

# The model

Kinetic model for the neutral atoms with Krook operators

$$\frac{\partial f_n}{\partial t} + \vec{v} \cdot \frac{\partial f_n}{\partial \vec{x}} = -\nu_{ion} f_n - \nu_{cx}(f_n - n_n \Phi_i) + \nu_{rec} n_i \Phi_i \quad (1)$$

$$\begin{aligned} \nu_{ion} &= n_e \langle \mathbf{v}_e \sigma_{ion}(\mathbf{v}_e) \rangle, & \nu_{cx} &= n_i \langle \mathbf{v}_{rel} \sigma_{cx}(\mathbf{v}_{rel}) \rangle \\ \nu_{rec} &= n_e \langle \mathbf{v}_e \sigma_{rec}(\mathbf{v}_e) \rangle, & \Phi_i &= f_i / n_i \end{aligned}$$

Boundary conditions

( $v_{\perp}$  in respect to the surface;  $\theta$  between  $\vec{v}$  and normal vector to the surface)

$$\int \vec{d}\mathbf{v} \, v_{\perp} f_n(\vec{x}_w, \vec{v}) + u_{i\perp} n_i = 0 \quad (2)$$

$$f_n(\vec{x}_w, \vec{v}) \propto \cos(\theta) e^{mv^2/2T_w} \quad \text{for } v_{\perp} > 0$$



# The model

Kinetic model for the neutral atoms with Krook operators

$$\frac{\partial f_n}{\partial t} + \vec{v} \cdot \frac{\partial f_n}{\partial \vec{x}} = -v_{ion} f_n - v_{cx}(f_n - n_n \Phi_i) + v_{rec} n_i \Phi_i \quad (1)$$

$$v_{ion} = n_e \langle v_e \sigma_{ion}(v_e) \rangle, \quad v_{cx} = n_i \langle v_{rel} \sigma_{cx}(v_{rel}) \rangle$$

$$v_{rec} = n_e \langle v_e \sigma_{rec}(v_e) \rangle, \quad \Phi_i = f_i / n_i$$

Boundary conditions

( $v_{\perp}$  in respect to the surface;  $\theta$  between  $\vec{v}$  and normal vector to the surface)

$$\int d\vec{v} v_{\perp} f_n(\vec{x}_w, \vec{v}) + u_{i\perp} n_i = 0 \quad (2)$$

$$f_n(\vec{x}_w, \vec{v}) \propto \cos(\theta) e^{mv^2/2T_w} \quad \text{for } v_{\perp} > 0$$

# The model

Kinetic model for the neutral atoms with Krook operators

$$\frac{\partial f_n}{\partial t} + \vec{v} \cdot \frac{\partial f_n}{\partial \vec{x}} = -v_{ion} f_n - v_{cx}(f_n - n_n \Phi_i) + v_{rec} n_i \Phi_i \quad (1)$$

$$v_{ion} = n_e \langle v_e \sigma_{ion}(v_e) \rangle, \quad v_{cx} = n_i \langle v_{rel} \sigma_{cx}(v_{rel}) \rangle$$

$$v_{rec} = n_e \langle v_e \sigma_{rec}(v_e) \rangle, \quad \Phi_i = f_i / n_i$$

Boundary conditions

( $v_{\perp}$  in respect to the surface;  $\theta$  between  $\vec{v}$  and normal vector to the surface)

$$\int d\vec{v} v_{\perp} f_n(\vec{x}_w, \vec{v}) + u_{i\perp} n_i = 0 \quad (2)$$

$$f_n(\vec{x}_w, \vec{v}) \propto \cos(\theta) e^{mv^2/2T_w} \quad \text{for } v_{\perp} > 0$$

# The model

Kinetic model for the neutral atoms with Krook operators

$$\frac{\partial f_n}{\partial t} + \vec{v} \cdot \frac{\partial f_n}{\partial \vec{x}} = -v_{ion} f_n - v_{cx}(f_n - n_n \Phi_i) + v_{rec} n_i \Phi_i \quad (1)$$

$$v_{ion} = n_e \langle v_e \sigma_{ion}(v_e) \rangle, \quad v_{cx} = n_i \langle v_{rel} \sigma_{cx}(v_{rel}) \rangle$$

$$v_{rec} = n_e \langle v_e \sigma_{rec}(v_e) \rangle, \quad \Phi_i = f_i / n_i$$

Boundary conditions

( $v_{\perp}$  in respect to the surface;  $\theta$  between  $\vec{v}$  and normal vector to the surface)

$$\int d\vec{v} v_{\perp} f_n(\vec{x}_w, \vec{v}) + u_{i\perp} n_i = 0 \quad (2)$$

$$f_n(\vec{x}_w, \vec{v}) \propto \cos(\theta) e^{mv^2/2T_w} \quad \text{for } v_{\perp} > 0$$

# The model in steady state

Steady state,  $\frac{\partial f_n}{\partial t} = 0$ , first approach

- ▶ Valid if  $\tau_{neutral\ losses} < \tau_{turbulence}$
- ▶ e.g.  $T_e = 20\text{eV}$ ,  $n_0 = 5 \cdot 10^{19}\text{m}^{-3}$

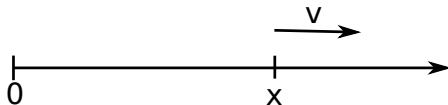
$$\tau_{loss} \approx v_{eff}^{-1} \approx 5 \cdot 10^{-7}\text{s}$$

$$\tau_{turbulence} \approx \sqrt{R_0 L_p} / c_{s0} \approx 2 \cdot 10^{-6}\text{s}$$

- ▶ Otherwise: time dependent model

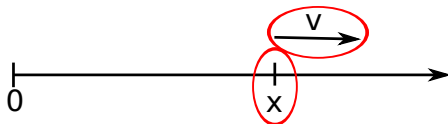
# The method of characteristics

Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$



# The method of characteristics

Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$

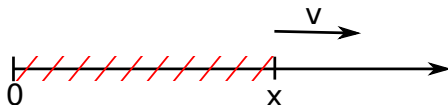


$$f_n(x, v)$$

(3)

# The method of characteristics

Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$

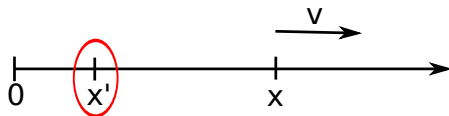


$$f_n(x, v) = \int_0^x dx'$$

(3)

# The method of characteristics

Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$

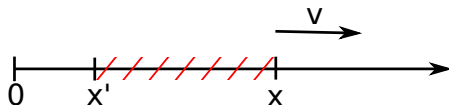


$$f_n(x, v) = \int_0^x dx' \frac{v_{cx}(x') n_n(x') \Phi_i(x', v)}{v} \quad (3)$$



# The method of characteristics

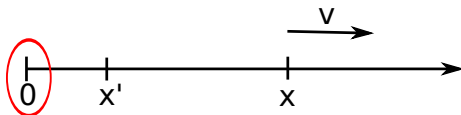
Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$



$$f_n(x, v) = \int_0^x dx' \frac{v_{cx}(x') n_n(x') \Phi_i(x', v)}{v} e^{\frac{1}{v} \int_{x'}^x dx'' (v_{cx}(x'') + v_{ion}(x''))} \quad (3)$$

# The method of characteristics

Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$

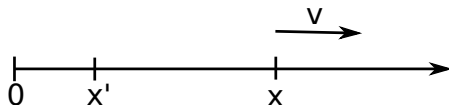


$$f_n(x, v) = \int_0^x dx' \frac{v_{cx}(x') n_n(x') \Phi_i(x', v)}{v} e^{\frac{1}{v} \int_{x'}^x dx'' (v_{cx}(x'') + v_{ion}(x''))} \quad (3)$$

$$+ f_w(v) e^{\frac{1}{v} \int_0^x dx'' (v_{cx}(x'') + v_{ion}(x''))}$$

# The method of characteristics

Example in 1D, no recombination,  $v > 0$  and a wall at  $x = 0$



$$f_n(x, v) = \int_0^x dx' \frac{v_{cx}(x') n_n(x') \Phi_i(x', v)}{v} e^{\frac{1}{v} \int_{x'}^x dx'' (v_{cx}(x'') + v_{ion}(x''))} \quad (3)$$

$$+ f_w(v) e^{\frac{1}{v} \int_0^x dx'' (v_{cx}(x'') + v_{ion}(x''))}$$

# An equation for the density distribution

By imposing  $\int dv f_n = n_n$

$$n_n(x) = \int_0^x dx' n_n(x') \int_0^\infty dv \frac{v_{cx}(x') \Phi_j(x', v)}{v} e^{\frac{v_{eff}(x-x')}{v}} \quad (4)$$

+ contribution by  $v < 0$   
+  $n_w(x)$

we get an integral equation for  $n_n(x)$ .

# Discretized equations

Discretize the spatial direction in  $x \rightarrow x_i$  and  $n_n(x_i) \rightarrow n_n^i$

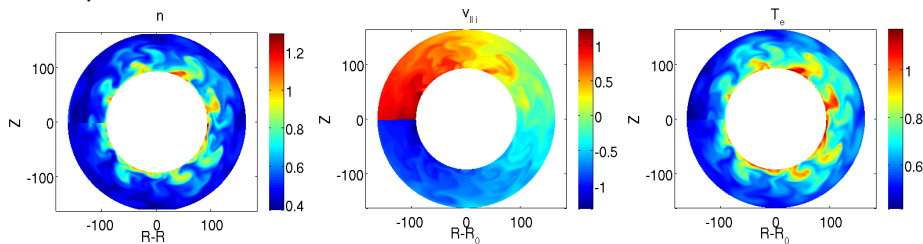
$$n_n^i = \sum_{j \leq i} n_n^j \int_0^\infty dv \frac{v_{cx}(x_j) \Phi_i(x_j, v)}{v} e^{\frac{v_{eff}(x_i - x_j)}{v}} \quad (5)$$

+ *contribution by  $v < 0$*   
+  $n_w(x_i)$

System of linear equations, solved with standard methods

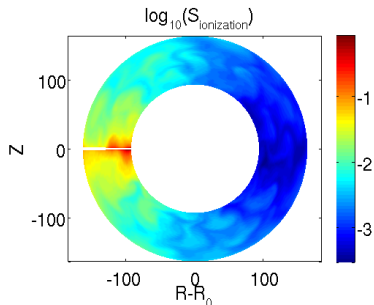
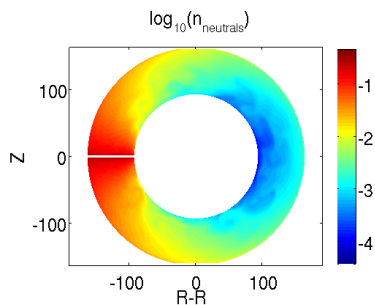
# Self-consistent simulation with GBS and neutrals model

Plasma: snapshot of density, parallel ion velocity and electron temperature from GBS



# Self-consistent simulation with GBS and neutrals model

Neutral density distribution and ionization rate ( $n_n n_e \langle v_e \sigma_{ion} \rangle$ ) from the simple neutral model,  $n_0 = 5 \cdot 10^{19} \text{m}^{-3}$ ,  $T_0 = 10 \text{eV}$



# Conclusions

- ▶ Developement of a neutral model for GBS
- ▶ Kinetic equation with Krook operators for ion, rec and cx
- ▶ 2D evaluation of the neutral density shows its proximity to the limiter



# Outlook

- ▶ Find the transition from convection to conduction limited regime
- ▶ Relax the steady state assumption
- ▶ Adding density source, drag force, and temperature source/sink to the equations in GBS
- ▶ Move towards detachment
  - ▶ Divertor geometry
  - ▶ Radiative detachment in limited geometry
  - ▶ Mimic the geometry of one divertor leg